

HYDRODYNAMIC PRESSURE IN HUMAN HIP JOINT

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Abstract

This paper presents modelling and simulations for synovial fluid flow occurring in gap between two co-operating bone surfaces in human hip joint. The present consideration gives an analysis of solutions of systems of non-linear, partial, differential equations for synovial fluid flow in human joint gap. In the hip joint the spherical rotation bone and the pelvis bone create a spherical gap. In this gap between two co-operating bones synovial fluid flows, see Dowson [1], Maquet [4], Mow [5], [6], [7], [8], [9], [10], Wierzcholski [12], [13], Pytko et.al. [11]. The flow of this fluid is caused by the motion of the bone head. The theoretical considerations of the synovial non-Newtonian fluid flow in thin joint gap, taking into account boundary layer simplifications, have practical applications in theory of lubrication in medicine, see Dowson et.al. [1], [2], Gruca [3]. The consideration given in the present section enables to find synovial fluid flow parameters and carrying capacity forces in human hip joint gap between two co-operating cartilage surfaces.

Human hip joint and gap height for deformable cartilage and pressure distribution in normal and pathological spherical hip joint gap for hydrodynamic lubrication caused by rotation, cases of pressure distribution in pathological spherical hip joint gap for hydrodynamic lubrication caused by rotation, lubrication surface, area=20.38 cm² are presented in the paper.

Keywords: normal, pathological hip, capacity, compressive stresses

CIŚNIENIE HYDRODYNAMICZNE W BIODROWYM STAWIE CZŁOWIEKA

Streszczenie

Praca przedstawia obliczenia numeryczne rozkładu ciśnienia hydrodynamicznego na sferycznej głowie kostnej stawu biodrowego wykonującego ruch obrotowy. Uwzględniona została deformacja powierzchni chrząstki stawowej oraz brane są pod uwagę własności nie Newtonowskie cieczy synowialnej poprzez zmiany lepkości dynamicznej wywołane zmianami prędkości deformacji. Ponadto uwzględnia się zarówno przypadek zdrowego jak i chorego stawu biodrowego. W tym stawie między dwoma współpracującą kości przepływa płyn synowialny, patrz Dowson [1], Maquet [4], Skoście [5], [6], [7], [8], [9], [10], Wierzcholski [12], [13], Pytko et.al. [11]. Przepływ tego płynu jest spowodowany przez ruch główki kości. Teoretyczne rozważania synowialnego nie-newtonowskiego przepływu płynu w cienkiej w cienkiej szczelinie stawu biorąc pod uwagę uproszczenia warstwy granicznej, mają zastosowania praktyczne teorii smarowania w medycynie, patrz Dowson et.al. [1], [2], Gruca [3]. Rozważanie podane w obecna artykule umożliwią znalezienie parametrów strumienia synowialnego płynu i zdolności do przenoszenia sił w szczelinie stawu biodrowego człowieka między dwoma współpracującymi powierzchniami chrząstki.

Staw biodrowy człowieka oraz wysokość szczeliny dla odkształcalnej chrząstki oraz rozkład ciśnienia w szczelinie zdrowego i chorego stawu biodrowego przy hydrodynamicznym smarowaniu, przypadki rozkładów ciśnienia w chorej, sferycznej szczelinie biodra dla hydrodynamicznego, smarowania wywołanego ruchem obrotowym, przypadki rozkładów ciśnienia w zdrowej, sferycznej szczelinie biodra dla hydrodynamicznego smarowania wywołanego ruchem obrotowym są przedstawione w artykule.

Słowa kluczowe: Zdrowe i chore biodra, nośność, naprężenia ściskające

1. Introduction

We take into account equations of conservation of momentum and continuity equation for the non-stationary flow of synovial incompressible fluid between two rotational roughness cartilage surfaces in the spherical co-ordinates. Imposing the proper boundary conditions on the basic equations the pressure p we obtain from the following Reynolds equation:

$$\frac{\partial}{\partial \varphi} \left(\frac{\varepsilon^3}{\eta} \frac{\partial p}{\partial \varphi} \right) + R^2 \sin \left(\frac{\vartheta}{R} \right) \frac{\partial}{\partial \vartheta} \left[\frac{\varepsilon^3}{\eta} \frac{\partial p}{\partial \vartheta} \sin \left(\frac{\vartheta}{R} \right) \right] = 6\omega R^2 \frac{\partial \varepsilon}{\partial \varphi} \sin^2 \left(\frac{\vartheta}{R} \right), \quad (1)$$

in Ω region: $0 \leq \varphi \leq \pi$, $\pi R/8 \leq \vartheta \leq \pi R/2$. We denote: η -dynamic viscosity of synovial fluid, ω -angular velocity of bone head, R -radius of bone head.

2. Deformable gap height

The minimum gap height, see Dowson [1], for spherical hip joint can be obtained from the formula [2]:

$$\frac{\varepsilon_{\min}}{R} \equiv \sqrt[5]{2\pi} S_1^{0.4} \left(\frac{\omega R^2 \eta_p}{C} \right)^{0.6}, \quad (2)$$

$$S_1 \equiv \frac{C}{ER^2}, \quad \frac{1}{E} = \frac{1}{2} \left(\frac{1-\nu_1^2}{E_1} + \frac{1-\nu_2^2}{E_2} \right),$$

where:

E_1, E_2, ν_1, ν_2 – the elastic module and Poisson ratios for bone head and cartilage,
 C – load, and the quantities η, ω, R are as defined previously.

Relation for share rate $\Theta \approx \omega R / \varepsilon_{\min}$ can be written in the following form:

$$\frac{\omega R^2 \eta_p}{C} \equiv \frac{S_2}{S_1} \left(\frac{\eta_o - \eta_\infty}{\eta_o + \frac{\eta_o}{1 + S_3 \frac{R}{\varepsilon_{\min}}}} \right), \quad S_2 \equiv \frac{\omega R \eta_o}{ER}, \quad S_3 \equiv A\omega. \quad (3)$$

By combining equations (2), (3) we obtain the system of two equations for determination of two unknown quantities, namely the dynamic viscosity η_p of synovial fluid and the minimum value ε_{\min} of gap height, where elastic deformations of cartilage are taken into account. If we assume the following data:

$R=2.6 \times 10^{-2}$ m, $E=2 \times 10^5$ Pa, $\omega R=3 \times 10^{-1}$ m/s, $\eta_\infty=0.10$ Pas, $2\pi R/C=3 \times 10^{-4}$ m/N, $\eta_o/\eta_\infty \cong 1000$, $A=1.88$ s, $C=544.26$ N, then from Eqs. (2), (3) we obtain:

$\varepsilon_{\min}=0.0000208$ m= 20.88 μ m and $\eta_p=0.1036$ Pas. If we take in computations the following quantities:

$A=1.88$ s, $\eta_o=100.00$ Pas, $\eta_\infty=0.10$ Pas, $R=0.020$ m, $C=544$ N, $0.50 \text{ s}^{-1} \leq \omega \leq 10.00 \text{ s}^{-1}$, $2 \times 10^5 \text{ Pa} \leq E \leq 2 \times 10^7 \text{ Pa}$, then we obtain the minimum value of gap height in the interval:

$0.29 \mu\text{m} \leq \varepsilon_{\min} \leq 19.90 \mu\text{m}$.

3. The gap height and primary calculations

The gap height has the following form:

$$\varepsilon = \Delta\varepsilon_x \cos\varphi \sin\vartheta / R + \Delta\varepsilon_y \sin\varphi \sin\vartheta / R - \Delta\varepsilon_z \cos\vartheta / R - R + [(\Delta\varepsilon_x \cos\varphi \sin\vartheta / R + \Delta\varepsilon_y \sin\varphi \sin\vartheta / R - \Delta\varepsilon_z \cos\vartheta / R)^2 + (R + \varepsilon_{\min})(R + 2D + \varepsilon_{\min})]^{0.5} \quad (4)$$

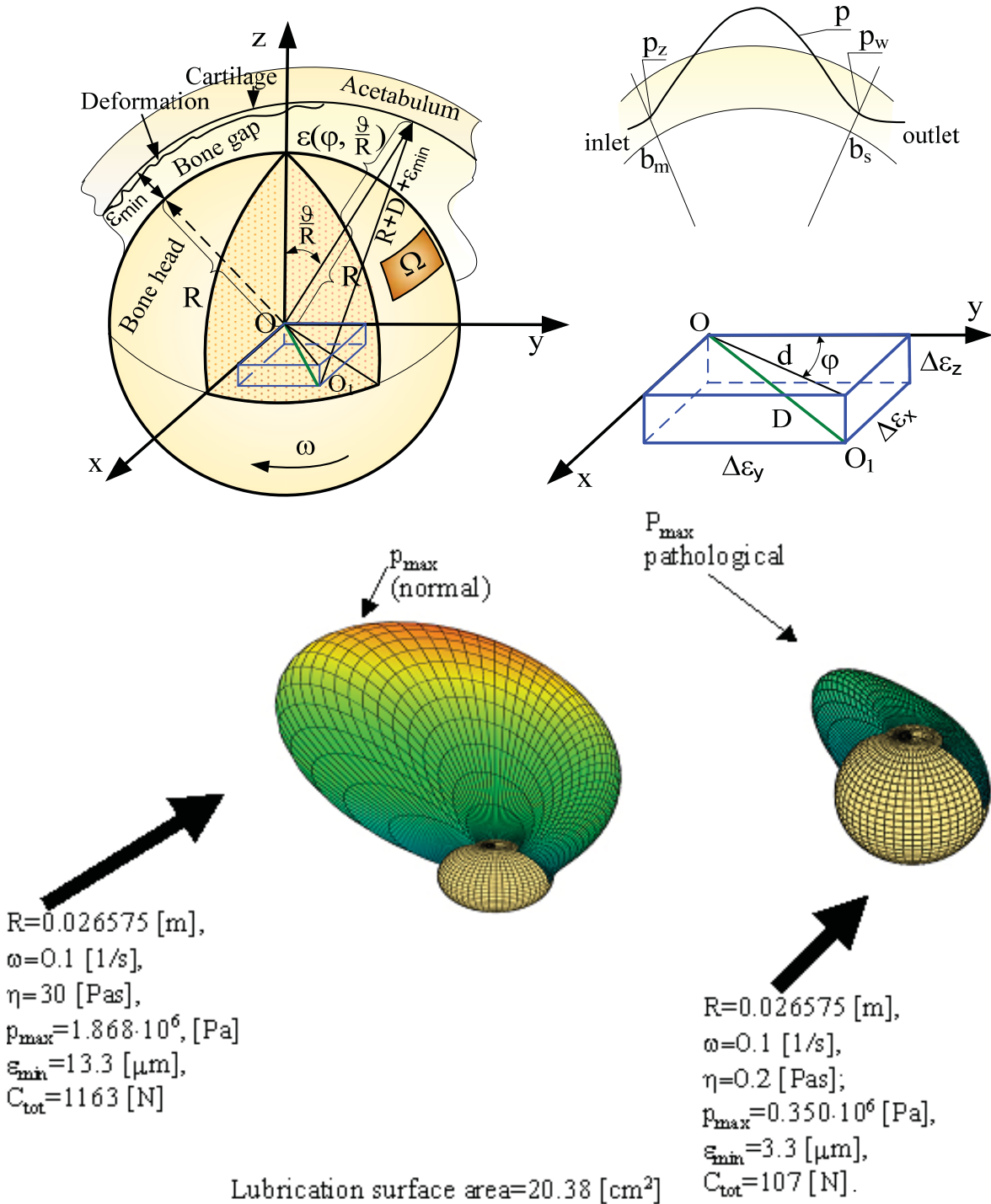


Fig. 1. Human hip joint and gap height for deformable cartilage and pressure distribution in normal and pathological spherical hip joint gap for hydrodynamic lubrication caused by rotation

Rys.1. Staw biodrowy człowieka oraz wysokość szczeliny dla odkształcalnej chrząstki oraz rozkład ciśnienia w szczelinie zdrowego i chorego stawu biodrowego przy hydrodynamicznym smarowaniu

The centre point of bone head (see Fig. 1) can be written in the following form:

$O_1(x - \Delta\varepsilon_x, y - \Delta\varepsilon_y, z + \Delta\varepsilon_z)$, while D is the distance between the centre of bone head and the acetabulum (sleeve) centre. We solve equation (1) for the region $\Omega(\varphi, \vartheta)$ resting on bone head and

indicated in Fig. 1. We assume atmospheric pressure on a boundary of the region $\Omega(\alpha_1, \alpha_3)$. The centre point of bone head (see Fig. 1) can be written in the following form: $O_1(x-\Delta\epsilon_x, y-\Delta\epsilon_y, z+\Delta\epsilon_z)$, while D is the distance between the centre of bone head and the acetabulum (sleeve) centre. We solve equation (1) for the region $\Omega(\varphi, \vartheta)$ resting on bone head and indicated in Fig. 1. We assume atmospheric pressure on a boundary of the region $\Omega(\alpha_1, \alpha_3)$. For this region we calculate also the capacity values. In numerical calculations we assume the following values for joint gap: $\Delta\epsilon_x=5 \mu\text{m}$, $\Delta\epsilon_y=5 \mu\text{m}$, $\Delta\epsilon_z=+5 \mu\text{m}$; the radius of bone head $R=0.026575 \text{ m}$. Now we calculate pressure distributions in normal and pathological hip joint.

For the angular velocity of bone head $\omega=0.1 \text{ s}^{-1}$, and the average value of synovial fluid dynamic viscosity $\eta=30.00 \text{ Pas}$ for normal joint and $\eta=0.20 \text{ Pas}$ for pathological joint, we obtain the smallest gap height $\epsilon_{\min}=13.3 \mu\text{m}$ for normal joint and $\epsilon_{\min}=3.3 \mu\text{m}$ for pathological joint.

The hydrodynamic pressure p has its maximum value equal to $18.68 \times 10^5 \text{ N/m}^2 \approx 18.30$ at for normal joint and $3.50 \times 10^5 \text{ N/m}^2$ for pathological joint, and the capacity $C_{\text{tot}}=1163 \text{ N}$ for normal joint and the capacity $C_{\text{tot}}=107 \text{ N}$ for pathological joint, (C_{tot} -total capacity, ϵ -gap height) see Fig. 1.

For the angular velocity of bone head $\omega=1.0 \text{ s}^{-1}$ and the average value of synovial fluid dynamic viscosity $\eta_o=2.00 \text{ Pas}$ for normal joint and $\eta_o=0.05 \text{ Pas}$ for pathological joint, we obtain the smallest gap height $\epsilon_{\min}=13.3 \mu\text{m}$ for normal joint and $\epsilon_{\min}=3.3 \mu\text{m}$ for pathological joint.

The hydrodynamic pressure p has its maximum value equal to $12.79 \times 10^5 \text{ N/m}^2 \approx 12.79$ at for normal joint and $7.26 \times 10^5 \text{ Pa}$ for pathological joint, and the capacity $C_{\text{tot}}=775 \text{ N}$ for normal joint and the capacity $C_{\text{tot}}=268 \text{ N}$ for pathological joint. Lubrication surface area has the value $\pi R^2 \cos \pi / 8 \approx 20.38 \text{ cm}^2$.

Numerical calculations are performed by means of Mathcad 12 Professional program with the help of the Method of Finite Differences. This method satisfies the requirement of stability of numerical solutions for the partial differential equation (1).

4. Numerical calculations

Now we calculate pressure distribution in pathological joint (see Fig. 2.). For the angular velocity of bone head:

$\omega=10.0 \text{ s}^{-1}$, $\omega=7.5 \text{ s}^{-1}$, $\omega=5.0 \text{ s}^{-1}$, $\omega=2.5 \text{ s}^{-1}$, $\omega=1.0 \text{ s}^{-1}$, $\omega=0.75 \text{ s}^{-1}$ and the average value of synovial fluid dynamic viscosity:

$\eta_o=0.0584 \text{ Pas}$, $\eta_o=0.0656 \text{ Pas}$, $\eta_o=0.074 \text{ Pas}$, $\eta_o=0.100 \text{ Pas}$, $\eta_o=0.1259 \text{ Pas}$, $\eta_o=0.1359 \text{ Pas}$, we obtain the following maximum values of hydrodynamic pressure:

$p_{\max}=1.383 \text{ at}$, $p_{\max}=1.322 \text{ at}$, $p_{\max}=1.242 \text{ at}$, $p_{\max}=1.164 \text{ at}$, $p_{\max}=1.082 \text{ at}$, $p_{\max}=1.067 \text{ at}$, and the carrying capacity:

$C_{\text{tot}}=25.1607 \text{ N}$, $C_{\text{tot}}=21.197 \text{ N}$, $C_{\text{tot}}=15.941 \text{ N}$, $C_{\text{tot}}=10.771 \text{ N}$, $C_{\text{tot}}=5.4242 \text{ N}$, $C_{\text{tot}}=4.3913 \text{ N}$, respectively.

Now we calculate pressure distribution in normal joint (see Fig. 3). For the angular velocity of bone head:

$\omega=10.0 \text{ s}^{-1}$; $\omega=7.5 \text{ s}^{-1}$, $\omega=5.0 \text{ s}^{-1}$, $\omega=2.5 \text{ s}^{-1}$, $\omega=1.0 \text{ s}^{-1}$, $\omega=0.75 \text{ s}^{-1}$ and the average value of synovial fluid dynamic viscosity:

$\eta_o=6.813 \text{ Pas}$, $\eta_o=9.261 \text{ Pas}$, $\eta_o=13.082 \text{ Pas}$, $\eta_o=23.163 \text{ Pas}$, $\eta_o=42.987 \text{ Pas}$, $\eta_o=50.118 \text{ Pas}$, we obtain the following maximum hydrodynamic pressure values:

$p_{\max}=45.63 \times 10^5 \text{ Pa} \approx 45.63 \text{ at}$, $p_{\max}=46.50 \times 10^5 \text{ Pa} \approx 46.50 \text{ at}$, $p_{\max}=43.75 \times 10^5 \text{ Pa} \approx 43.75 \text{ at}$, $p_{\max}=39.10 \times 10^5 \text{ Pa} \approx 39.10 \text{ at}$, $p_{\max}=29.16 \times 10^5 \text{ Pa} \approx 29.16 \text{ at}$, $p_{\max}=25.62 \times 10^5 \text{ Pa} \approx 25.62 \text{ at}$,

and the carrying capacity: $C_{\text{tot}}=2935.3 \text{ N}$, $C_{\text{tot}}=2992.5 \text{ N}$, $C_{\text{tot}}=2818.1 \text{ N}$, $C_{\text{tot}}=2505.6 \text{ N}$, $C_{\text{tot}}=1852.0 \text{ N}$, $C_{\text{tot}}=1619.4 \text{ N}$, respectively.

Lubrication surface area has the value $R^2 \cos \pi / 8 \approx 20.38 \text{ cm}^2$.

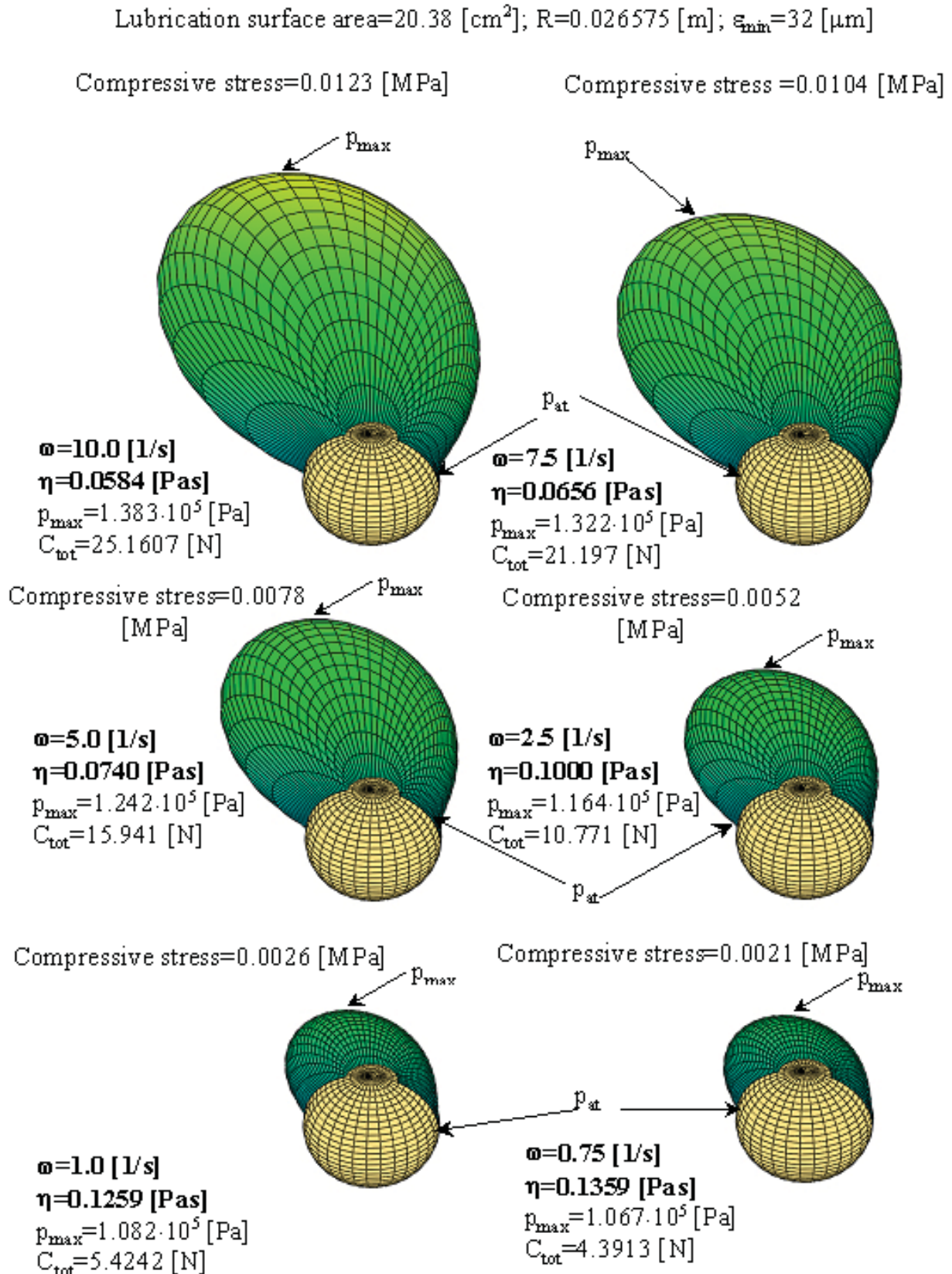


Fig. 2. Cases of pressure distribution in pathological spherical hip joint gap for hydrodynamic lubrication caused by rotation, lubrication surface, area=20.38 cm²
 Rys. 2. Przypadki rozkładów ciśnienia w chorej, sferycznej szczelinie biodra dla hydrodynamicznego smarowania wywołanego ruchem obrotowym

Lubrication surface area=20.38 [cm²]; R=0.026575 [m]; ε_{\min} =32 [μm]

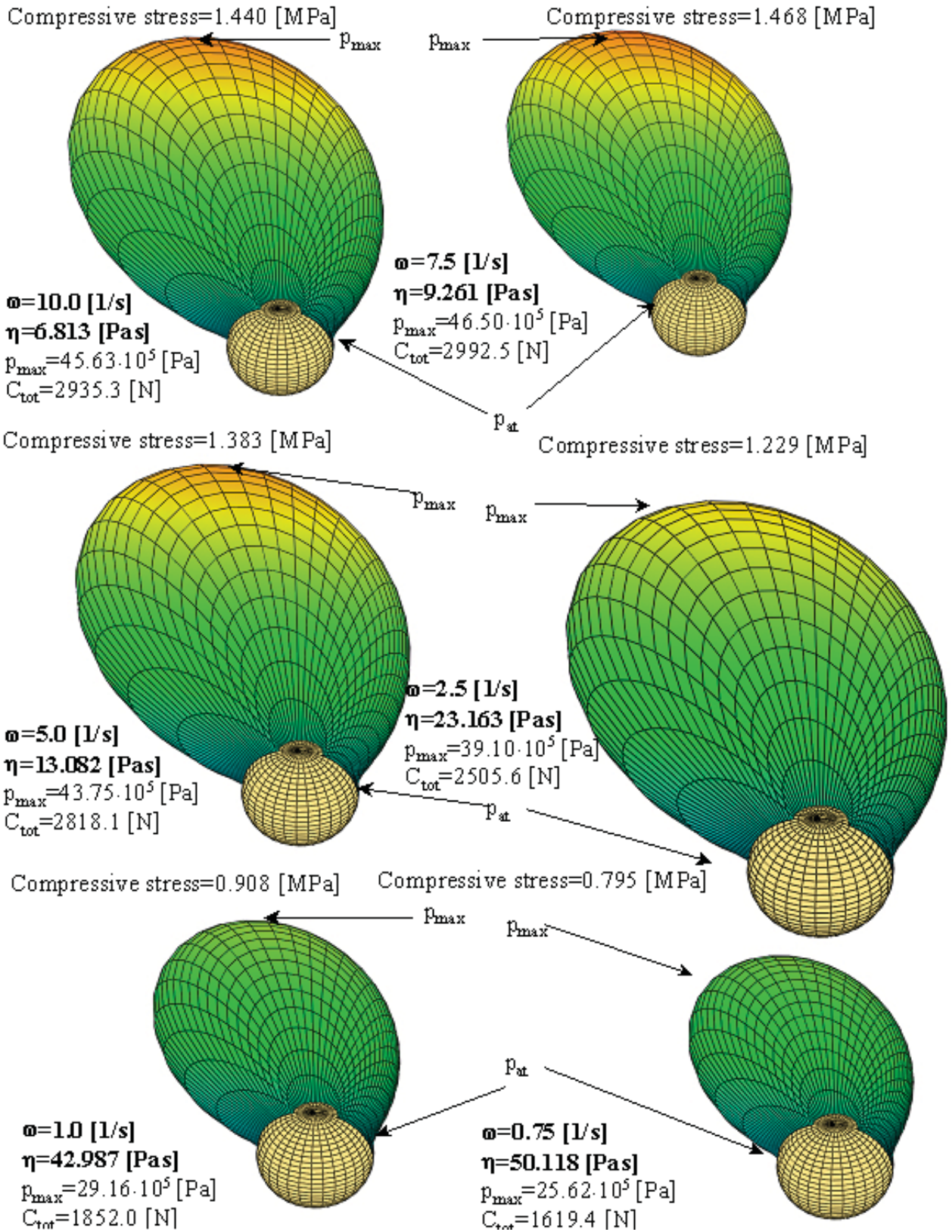


Fig. 3. Cases of pressure distribution in normal spherical hip joint gap for hydrodynamic lubrication caused by rotation, lubrication surface area =20.38 cm²

Rys. 3. Przypadki rozkładów ciśnienia w zdrowej, sferycznej szczelinie biodra dla hydrodynamicznego smarowania wywołanego ruchem obrotowym

5. Final comments

For the above given capacities in normal joint we obtain the following compressive stresses: $\tau_s = 2935.3 \text{ N} / 20.38 \text{ cm}^2 = 1.440 \text{ N/mm}^2 = 1.440 \text{ MN/m}^2$, $\tau_s = 1.468 \text{ MN/m}^2$, $\tau_s = 1.383 \text{ MN/m}^2$, $\tau_s = 1.229 \text{ MN/m}^2$, $\tau_s = 0.908 \text{ MN/m}^2$, $\tau_s = 0.795 \text{ MN/m}^2$.

In pathological joint, the compressive stresses are as follows: $\tau_s = 25.16 \text{ N} / 20.38 \text{ cm}^2 = 0.0123 \text{ N/mm}^2 = 0.0123 \text{ MN/m}^2$, $\tau_s = 0.0104 \text{ MN/m}^2$, These stresses are smaller than the compressive strength 21 MN/m^2 of normal human cartilage. Compressive strength of pathological human cartilage has the value of about 1 MPa [2], [3].

The present paper shows the method of determination of approximate solutions of the partial non-linear differential equations of non-Newtonian, asymmetrical synovial fluid flow in the thin gap occurring in human joint in curvilinear, orthogonal co-ordinates.

The presented method enables to obtain solutions in the form of Taylor series with increasing powers of the small parameter A obtained in experimental way for synovial fluid.

In the particular case of the symmetrical flow we can, by virtue of the presented theory, find analytical solutions in a simple form. The percentage corrections of pressure caused by the non-Newtonian properties of synovial fluid are examined numerically. For large shear rates within the range: $100 \text{ s}^{-1} \leq \Theta \leq 1000 \text{ s}^{-1}$, the viscosity of synovial fluid is small and has values: $10^{-1} \text{ Pas} \leq \eta \leq 1 \text{ Pas}$. In this case we obtain small pressure changes from 2% to 4%. For small shear rates within the range: $10^{-1} \text{ s}^{-1} \leq \Theta \leq 10 \text{ s}^{-1}$, when viscosity of synovial liquid is large, i.e. within the range: $10 \text{ Pas} \leq \eta \leq 100 \text{ Pas}$, from Eq. (5) we obtain pressure changes from 7% to 15%.

- Note 1: If Young Modulus of cartilage and bone decreases then the least value of joint gap height increases.

- Note 2: If Young Modulus of cartilage decreases then viscosity of synovial fluid increases.

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5. References

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